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Contrasting interannual and multidecadal NAO variability

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Abstract Decadal and longer timescale variability in the winter North Atlantic Oscillation (NAO) has considerable impact on regional climate, yet it remains unclear what fraction of this variability is potentially predictable. This study takes a new approach to this question by demonstrating clear physical differences between NAO variability on interannual-decadal (<30 year) and multidecadal (>30 year) timescales. It is shown that on the shorter timescale the NAO is dominated by variations in the latitude of the North Atlantic jet and storm track, whereas on the longer timescale it represents changes in their strength instead. NAO variability on the two timescales is associated with different dynamical behaviour in terms of eddy-mean flow interaction, Rossby wave breaking and blocking. The two timescales also exhibit different regional impacts on temperature and precipitation and different relationships to sea surface temperatures. These results are derived from linear regression analysis of the Twentieth Century and NCEP-NCAR reanalyses and of a high-resolution HiGEM General Circulation Model control simulation, with additional analysis of a long sea level pressure reconstruction. Evidence is presented for an influence of the ocean circulation on the longer timescale variability of the NAO, which is particularly clear in the model data. As well as providing new evidence of potential predictability, these findings are shown to have implications for the reconstruction and interpretation of long climate records.

Keywords North Atlantic Oscillation · Jet variability · Atmosphere-ocean interaction · Climate reconstructions

1 Introduction

As the leading pattern of atmospheric circulation variability over the North Atlantic, the North Atlantic Oscillation¹ (NAO) has a strong influence on surface climate across the Atlantic basin and beyond (Thompson and Wallace 2001). Interest in the NAO has been partly motivated by the prominence of its decadal-scale variability in winter (Stephenson *et al.* 2000). The increase of the winter NAO index from the 1960s to the 1990s gained particular attention (Hurrell 1995) but decadal variability is also evident in longer records of the NAO (Pinto and Raible 2012). In the likely absence of atmospheric memory from one winter season to the next, influences from other components of the climate system may have played a role. Evidence has been provided of possible influences such as the extratropical (Rodwell *et al.* 1999; Czaja and Frankignoul 1999; Mosedale *et al.* 2006; Gastineau and Frankignoul 2012) or tropical oceans (Hoerling *et al.* 2001; Selten *et al.* 2004; Greatbatch *et al.* 2012), the sea-ice (Deser *et al.* 2004; Bader *et al.* 2011) and also forcings acting via the stratosphere, such as changes in stratospheric water vapour (Joshi *et al.* 2006) or solar variability (Ineson *et al.* 2011). The associated potential for predictability of NAO variability continues to drive research in this area (Folland *et al.* 2012).

Much recent work has focused on shorter, intraseasonal timescales in attempts to understand the atmospheric dynamics underlying NAO variability,

¹ Or equivalently, the Arctic Oscillation or Northern Annular Mode (Feldstein and Franzke 2006).

following Feldstein (2003) and Benedict *et al.* (2004). It is clear that some of the NAO variability on decadal timescales could arise from so-called climate noise, in which seasonal sampling of the strong intraseasonal variability can lead to apparent power on interannual and longer timescales (Wunsch 1999; Feldstein 2000b; Schneider *et al.* 2003; Raible *et al.* 2005). Various statistical methods have been applied to estimate the fraction of variance on interannual and longer timescales which could be explained simply as climate noise. However, these methods differ widely in their findings (Feldstein 2000a; Keeley *et al.* 2009; Franzke and Woollings 2011), so that the statistical significance of low-frequency NAO variability, and hence the potential for seasonal-decadal predictability is still unclear. In this paper we take a different and complementary approach, by searching for physical differences between NAO variability on short and long timescales.

The NAO is essentially a description of the preferred structure of variability in the North Atlantic eddy-driven jet stream (Thompson *et al.* 2002). This deep tropospheric jet represents the net effect of westerly wind forcing by the transient atmospheric eddies (Li and Wettstein 2012), and variations in its strength and position affect regional temperatures and precipitation via variations in the prevailing westerly winds and associated storm tracks. The NAO is usually defined via patterns in surface pressure or geopotential height, using methods such as principal component analysis. Physical quantities such as the latitude and speed of the jet are generally not separable by these methods (Monahan *et al.* 2009), and the NAO reflects variations in both of these quan-

73 titles (Woollings *et al.* 2010). When height fields are linearly regressed onto
74 time series of the jet latitude and speed, both of the resulting spatial pat-
75 terns project onto the NAO pattern (Woollings and Blackburn 2012). Despite
76 this, the jet latitude and speed are clearly distinct, having different annual cy-
77 cles, power spectra and interannual variability (Woollings and Blackburn 2012;
78 Woollings *et al.* 2014). This suggests that variations in jet latitude and speed
79 have different physical mechanisms and drivers, and yet they are combined in
80 standard NAO analyses.

81 Here we highlight the contrasting nature of the jet variability associated
82 with the NAO on two different timescales, namely multidecadal and interannual-
83 decadal, with periods greater or less than 30 years respectively. This study
84 is related to other approaches which focus on the non-stationarity of the
85 NAO pattern over time (Jung *et al.* 2003; Lu and Greatbatch 2002; Raible
86 *et al.* 2006; Wang *et al.* 2012; Moore *et al.* 2013), or multi-decadal changes
87 in regime activity (Casty *et al.* 2005; Franzke *et al.* 2011). Other studies have
88 highlighted non-stationary relationships between the NAO and regional im-
89 pacts on temperature (Pozo-Vázquez *et al.* 2001; Haylock *et al.* 2007; Comas-
90 Bru and McDermott 2013), precipitation (Vicente-Serrano and López-Moreno
91 2008; Raible *et al.* 2014) and storm activity (Luo *et al.* 2011; Lee *et al.*
92 2012), and the timescale dependence shown here may help in interpreting
93 this non-stationarity. Finally, there is evidence for distinct patterns of ocean-
94 atmosphere variability on decadal/multi-decadal timescales in observations
95 and models (Deser and Blackmon 1993; Delworth and Mann 2000; Sutton and

Hodson 2003; Shaffrey and Sutton 2006), with non-stationarity or frequency-dependence in the relationship between the NAO and sea surface temperatures (Raible *et al.* 2001; Walter and Graf 2002; Raible *et al.* 2005; Alvarez-Garcia *et al.* 2008). Hence we also examine the NAO-SST relationship on the two timescales as a preliminary study of the associated ocean-atmosphere interaction.

2 Methods

In this study we focus on the variability in wintertime (DJF) mean data from atmospheric reanalyses. We use both the NCEP-NCAR reanalysis (Kalnay and coauthors 1996) and the Twentieth Century reanalysis (20CR) (Compo *et al.* 2011). The latter uses mean sea level pressure (MSLP) observations only and takes an ensemble approach to quantify uncertainty, providing 56 estimates of atmospheric flow from 1871 to 2012. Unless otherwise stated, the analyses presented here were performed individually for each of these ensemble members and only averaged over the ensemble at the end of the analysis.

The NCEP-NCAR reanalysis is used over the period 1950-2012. Since this only provides 62 winters of data this record is short to examine the multi-decadal behaviour, but as will be shown results are qualitatively similar to those in 20CR. This reanalysis is used particularly to investigate transient features such as the storm track and the transient eddy fluxes. These quantities have been examined in 20CR but found to give unphysical results in the low-frequency regressions, in particular at high latitudes. This is likely due

to issues in the data-sparse period before 1920 (see e.g. Krueger *et al.* (2013) for a general discussion). The storm track is characterised using data filtered with a 2-6 day Chebyshev recursive filter to select only the synoptic timescales (Cappellini 1978).

We also make use of a 100 year present-day control simulation of the high-resolution coupled General Circulation Model HiGEM (Shaffrey *et al.* 2009). This has an atmospheric resolution of $0.833^\circ \times 1.25^\circ$ in longitude-latitude with 38 levels, and an ocean resolution of $\frac{1}{3}^\circ$ with 40 levels. This model shows improved simulation of the climatology and variability of North Atlantic climate compared to the standard resolution HadGEM1.2 (Shaffrey *et al.* 2009; Keeley *et al.* 2012; Hodson and Sutton 2012). Some limited transient diagnostics have been derived from the HiGEM simulation and these agree well with the results from the NCEP-NCAR reanalysis.

The NAO was defined in all datasets as the leading Empirical Orthogonal Function (EOF) and associated principal component time series of monthly mean wintertime (DJF) mean sea level pressure over the Atlantic sector (90°W - 30°E , 30 - 90°N). In the 20CR data the NAO was calculated separately in each ensemble member, and the resulting average spatial pattern is shown in Figure 5a of Woollings *et al.* (2014, W14 hereafter). The monthly NAO index was averaged up to seasonal mean values for analysis. As described in W14, indices of jet latitude and speed were derived using the zonal wind at 850 hPa. The method essentially averages the daily zonal wind over 0 - 60°W and smooths it using a 10-day low pass filter before locating the maximum wind speed

(Woollings *et al.* 2010). The resulting daily values of jet latitude and speed were averaged over each winter season to derive seasonal mean values.

To separate the different timescales we apply Empirical Mode Decomposition (EMD), as in Franzke and Woollings (2011), to the seasonal mean time series of the NAO and jet indices. This approach empirically decomposes a time series into Intrinsic Mode Functions (IMFs) of different average periods. See Franzke and Woollings (2011) for more description and an example of the method. Here we focus on two timescales: the interannual-decadal, formed by isolating the IMFs with average periods less than 30 years, and the multi-decadal, with IMF periods greater than 30 years. The sum of the two filtered time series is exactly equal to the full unfiltered series. These two timescales were chosen after experimentation to best represent the contrasting NAO behaviour (for example the 10-30 year band of timescales behave similarly to the 1-10 year band). Note that the general results presented here are reproduced using other filtering methods such as running means, but the EMD results are presented due to their smoothness and objectivity.

The general approach taken here is to linearly regress various fields onto the NAO time series at the two different timescales. After averaging the monthly data up to seasonal means, the NAO series is re-normalised so that the series of winter mean values has a mean of zero and a standard deviation of one. Maps therefore show the anomalies associated with one standard deviation of the full unfiltered winter NAO. As described below, this makes the magnitudes of the patterns on the two timescales comparable. However, it is important

to note that the long timescale anomaly patterns then have larger amplitude than is experienced in practise.

3 Jet Characteristics

We begin by comparing the NAO and jet indices from 20CR in Figure 1. The decadal variability of the NAO is clear, with high NAO values dominating in the early and late twentieth century, and low NAO values dominating in the middle of the century. In contrast, the jet latitude shows mostly interannual variability, and as shown by W14 it is the jet speed which exhibits the strongest decadal variability. W14 used a Monte Carlo statistical test to assess the probability that the observed variability in the decadal means of these jet series could arise from a white noise process. The results showed that this was quite plausible for the jet latitude ($p=0.19$) but very unlikely for the jet speed ($p=0.01$).

W14 also found that the jet latitude and speed series are uncorrelated ($r=-0.07$), yet both are related to the NAO. This is shown in Figure 2 which correlates these series with the NAO series on the different timescales obtained using the EMD filtering. Both 20CR and NCEP-NCAR results are plotted, with errorbars reflecting the uncertainty across the ensemble in 20CR. On timescales shorter than 30 years the NAO is dominated by variations in jet latitude. However, on the multidecadal timescale the reverse is true for 20CR at least; the jet speed is more highly correlated with the NAO. NCEP-NCAR shows high correlations for both jet speed and jet latitude on this timescale. If

the 20CR analysis is restricted to the time period of the NCEP-NCAR reanalysis, this gives very similar correlations to the NCEP-NCAR data (asterisks in Figure 2), suggesting that this difference is largely due to the short time period.

These correlations suggest a change in the nature of NAO variability on long timescales, with variations in jet speed becoming more important. This impression is confirmed in Figure 3 which shows the 850 hPa zonal wind anomalies associated with NAO variability on the two timescales. On the shorter timescale the wind anomalies generally straddle the mean jet, indicating a meridional shift, although the anomalies exhibit weaker meridional tilt and are focused downstream of the mean wind maximum. On the multidecadal timescale, however, the anomalies overlie the mean jet, indicating a clear increase in jet speed during positive NAO variations. The increase in speed is also shifted towards the eastern end of the jet, highlighting an extension of the jet towards central Europe. Similar patterns are seen in the NCEP-NCAR data, despite the difference in correlation on the long timescale in Figure 2. Although the multidecadal anomalies are weaker in NCEP-NCAR, the zonal wind is strengthened along the jet core as seen in 20CR. The same behaviour is also seen in the HiGEM model control simulation, which suggests the result is not a coincidence of the recent observed period. The similarity of the patterns across the three datasets adds considerable confidence to the result. These other datasets will be used in particular to analyse the storm activity and ocean-atmosphere interaction on the two timescales, since these

are two aspects which have proved problematic in the analysis of the 20CR data.

Following the jet analyses of W14, we performed a simple statistical test of the decadal NAO variability against a white noise hypothesis. For each of the 56 ensemble members, 1000 surrogate white noise NAO time series were generated with the same standard deviation, and then decadal means were calculated from these. The surrogate series were then used to determine the likelihood of the observed level of variability in decadal means occurring from the noise. This analysis showed that the decadal NAO variability in 20CR is very unlikely to occur in a white noise model ($p=0.01$). We then applied multiple linear regression to express the NAO as a linear combination of the jet indices (which explained 71% of the NAO variance). This enabled us to remove the influence of jet latitude and speed in turn and recalculate the likelihood of the resulting decadal NAO variability. Removing the contribution of jet latitude variations resulted in an NAO series which was still very unlikely in the noise model ($p=0.01$), but removing the contribution of jet speed gave a value of $p=0.19$, so that the resulting decadal NAO variability was no longer significantly different from that expected from white noise. This exercise shows that it is the variations in jet speed which are responsible for the elevated power of the NAO on decadal timescales.

It would clearly be beneficial to verify the contrasting NAO behaviour on the two timescales in a longer observational dataset. Several attempts have been made to reconstruct atmospheric flow fields beyond the last century using

instrumental records, with considerable success. We have analysed the Küttel *et al.* (2010) reconstruction of MSLP over Europe and the eastern North Atlantic back to 1750, which uses both terrestrial pressure and marine wind data. The result is that this dataset does not exhibit the distinct nature of multi-decadal variability shown in the other datasets. Our analysis (described in section 8) suggests that this may be at least partly an artefact of an assumption of stationarity in the method used to derive the reconstruction. Given the strong agreement between the other datasets, we conclude that the reconstruction likely underestimates the timescale dependence of NAO variability.

4 Regional Impacts

NAO variations are of particular interest because of their strong influence on regional surface climate. These connections have obvious societal impact and are also often used to reconstruct indices of the NAO back in time. Figure 4 shows the patterns of near surface air temperature and precipitation associated with the NAO. These impacts are notably different on the two timescales, especially over Europe. On interannual timescales this analysis gives the canonical patterns of a quadrupole in temperature anomalies and a north-south dipole in precipitation. On the decadal timescale, however, these patterns are shifted south, so that both temperature and precipitation anomalies are focused on western-central Europe. This southward shift is consistent with the role of the East Atlantic pattern (the second EOF) which also describes changes in jet speed and can be interpreted as acting to shift the NAO circulation pattern

north and south (Woollings *et al.* 2010). Figure 4 also shows strong differences in Arctic climate on the two timescales, though this should be treated with caution. If the analysis is restricted to the first 100 years of the period (1871-1970) then the Arctic signal is greatly reduced. This suggests that this signal arises from a correlation with the recent Arctic warming trend, which may be unrelated to the NAO.

Some consideration of the variance of the NAO on the different timescales should be taken in interpreting these impacts. This applies to all of the regression maps shown in this paper. As described above, the NAO time series was normalised and then split into the two components. The variance of the full time series is 1.0, but the two components have very different variances: 0.90 for the interannual-decadal series and only 0.08 for the multidecadal series. This means that anomalies of the magnitude of those in the multidecadal regressions are never realised in practise; the units of these regressions are changes per standard deviation of the full NAO series, while the multidecadal changes are much smaller than this. Despite this, the anomalies are of considerable importance when compared to the level of variability on this timescale. Over western Europe the multidecadal NAO regression accounts for up to 50% of the variance in decadal mean zonal wind, and similarly 30% of the temperature and 60% of the precipitation variance.

5 Hemispheric connections

In this section we investigate whether the NAO on the two timescales has different links to remote regions, in particular the Pacific. This is motivated by discussion over the hemispheric or regional nature of the NAO (Wallace 2000) and also by evidence that interaction between the Atlantic and Pacific sectors might be non-stationary or timescale dependent (Raible *et al.* 2001; Castanheira and Graf 2003; Pinto *et al.* 2011; Lee *et al.* 2012).

Figure 5 shows the MSLP associated with NAO variations on the two timescales. Despite the clear differences in jet behaviour, the MSLP patterns in the Atlantic are only subtly different, indicating that MSLP anomalies are hard to interpret in terms of jet characteristics. The most robust difference is an eastward shift of the equatorward Atlantic centre of action on the long timescale. The two reanalyses and the model are in good agreement over the structure of variability on the shorter timescale, including a weak centre of action in the eastern North Pacific. Climate models have historically overestimated the NAO teleconnection to the North Pacific (McHugh and Rogers 2005), but HiGEM appears to perform well in this regard (at least on the short timescale). On the longer timescale there is little agreement between the three datasets in the Pacific sector (though again the two reanalyses are similar if only the NCEP-NCAR period is used; not shown). This lack of agreement between datasets limits the confidence we can have in hemispheric connections on the long timescale.

To further investigate the Atlantic-Pacific links in the observations, we analyse the storm track variability in the NCEP-NCAR reanalysis. The upper level storm track is summarised by the mean of v'^2 at 200 hPa, which is regressed onto the NAO in Figure 6. In the Atlantic sector the contrast in storm track behaviour is very clear. The positive phase of the NAO is associated with a northward shift and extension of the storm track on interannual-decadal timescales, whereas it is associated with a strengthening of the storm track on multidecadal timescales. These features do extend upstream into the Pacific on both timescales. This is particularly clear on the multidecadal timescale, where a strong increase in storm activity is seen over the eastern North Pacific. This is consistent with the results of Lee *et al.* (2012) who found similar long term changes in the Atlantic and Pacific storm tracks over recent decades. Although weaker, the Pacific storm track signal on the shorter timescale is again consistent with the Atlantic flow, since it indicates a weakening on the southern side of the storm track. It is also consistent with studies which have noted a latitudinal shift of the Pacific storm track accompanying an Atlantic shift (Franzke *et al.* 2004; Strong and Magnusdottir 2008). Figure 6 also shows corresponding results from HiGEM. As in the MSLP analysis, there seems to be a Pacific-Atlantic storm track link on the shorter timescale which agrees well with that in the reanalysis. On the long timescale there is good agreement between the model and the reanalysis over the Atlantic sector, but not over the Pacific.

To summarise, there is evidence that NAO variability on both timescales has links to the Pacific sector. This is particularly clear on the shorter timescale, where both MSLP and storm tracks show good agreement between the different datasets. In contrast, confidence in Atlantic-Pacific links on the longer timescale is limited by the large differences between the three datasets.

6 Eddy-mean flow interaction

In this section we present further dynamical diagnostics of the atmospheric circulation differences on the two timescales. We use the NCEP-NCAR reanalysis for this analysis because of higher confidence in its transient fields. Figure 7 shows the vertical structure of the zonal wind anomalies along a section at 30°W . On both timescales the wind anomalies are equivalent barotropic with maxima in the upper troposphere. The differences between the two timescales seen at 850 hPa are clearly evident in the eddy-driven jet through the depth of the troposphere, and are not just surface features. This suggests that the wind anomalies are accompanied by changes in transient eddy driving of the zonal flow, as expected from the storm track changes shown in Figure 6. Figure 7 also shows an interesting contrast in subtropical jet variability. This is opposite to the eddy-driven jet, in that the subtropical jet strengthens and weakens on the short timescale but shifts meridionally on the long timescale.

Transient baroclinic eddies influence the large-scale flow via both heat and momentum fluxes. The top panels of Figure 8 show the lower tropospheric transient eddy heat fluxes ($v'T'$). As in the other fields there is a clear

change between the timescales from a largely shifting pattern of variability to a strengthening one. The transient eddy heat fluxes drive a residual overturning circulation with Coriolis torque acting to accelerate the westerly flow at the latitude of the maximum in $v'T'$. Figure 8 then shows that the changes in transient eddy heat fluxes act to support the zonal flow variations in each case, helping to shift the surface westerlies in the interannual-decadal variability and strengthen them in the multidecadal variability.

To summarise the effects of the transient eddy momentum fluxes, we follow Raible *et al.* (2010) in calculating $\mathbf{E} \cdot \mathbf{D}$ where $\mathbf{E} = ((v'^2 - u'^2)/2, -u'v')$ is similar to the E-vector of Hoskins *et al.* (1983) and $\mathbf{D} = (U_x - V_y, V_x + U_y)$ is the deformation vector of the time mean flow. Here u' and v' are the 2-6 day band-pass filtered wind components and U and V are the wind components averaged over the relevant winter season. This diagnostic describes the exchange of kinetic energy between the eddies and the background flow (Mak and Cai 1989). Regressions of $\mathbf{E} \cdot \mathbf{D}$ on the NAO are shown in the lower panels of Figure 8. The climatology of $\mathbf{E} \cdot \mathbf{D}$ features positive values over North America, implying that eddies grow there at the expense of the background state. Over the Atlantic Ocean the climatology is negative, showing that the eddies lose kinetic energy to the background state there. The regression patterns are again very different on the two timescales. The multidecadal regression shows a strengthening of the conversion from eddy to background state kinetic energy, consistent with increased eddy driving of the stronger jet stream. On the interannual-decadal timescale the pattern is more complicated. While the

region of maximum eddy forcing is shifted northward by the anomalies, the pattern also shows a meridional tightening of the eddy forcing over the ocean and a general strengthening downstream. The upstream part of this pattern may be related to the strengthening of the subtropical jet seen in Figure 7.

The effect of eddy forcing on the mean flow of the NAO is increasingly described with regard to the breaking of transient Rossby waves (Benedict *et al.* 2004; Franzke *et al.* 2004; Rivière and Orlanski 2007; Martius *et al.* 2007; Kunz *et al.* 2009; Archambault *et al.* 2010). Here we use the index of Barnes and Hartmann (2012) which identifies wave breaking via the latitudinal overturning of vorticity contours. The index outputs the centroid of the wave breaking event, counting each event once only (with a median lifetime of events of two days) and discriminates between cyclonic and anticyclonic breaking based on the morphology of the overturning region. Regressions of the occurrence of wave breaking on the NAO are shown in Figure 9. On the interannual-decadal timescale, the positive NAO is associated with a reduction in cyclonic wave breaking on the poleward flank of the jet (to the south of Greenland) and an increase in wave breaking of both types on the equatorward flank of the jet. There is also a decrease in anticyclonic breaking in the subtropics, suggesting that this region of wave breaking shifts north along with the jet. These patterns are consistent with the picture that Rossby wave breaking acts to decelerate the westerly winds locally, so that breaking on the equatorward side pushes the jet polewards and vice versa (Gabriel and Peters 2008).

On the multidecadal timescale the positive NAO is instead associated with increased wave breaking on both sides of the jet, which is consistent with the strengthening and extension of the jet. Such large-scale conditions are known to foster the occurrence of extreme windstorms over Western Europe (Hanley and Caballero 2012; Gómara *et al.* 2014). The strongest signals are increased cyclonic breaking to the north and increased anticyclonic breaking to the south, though the two types of breaking also show weaker increases on the opposite side of the jet. On both timescales the behaviour is therefore consistent with Strong and Magnusdottir (2008), in that the latitude of the breaking seems more important than its direction (e.g. the breaking on the equatorward side of the jet may be cyclonic as well as anticyclonic).

Finally in this section, we analyse the relationship between the NAO and blocking on both timescales. Blocking is a synoptic situation in which the westerly winds and storm tracks are blocked by a persistent, usually anticyclonic, flow anomaly. Blocking is itself related to wave-breaking (Pelly and Hoskins 2003; Altenhoff *et al.* 2008), though the requirements of spatial scale and persistence separate it from more transient wave breaking (Masato *et al.* 2009). Here we define blocking using the index of Scherrer *et al.* (2006) which is a two-dimensional extension of the classical Tibaldi and Molteni (1990) index. A blocking pattern is identified at a point if 1) the meridional 500 hPa geopotential height gradient is reversed and 2) the flow is westerly to the north of the point, with a height gradient stronger than 10 m per degree of latitude. A 5-day persistence criterion is then applied at each gridpoint.

Figure 10 shows the regressions of blocking activity on the NAO. On the interannual-decadal timescale, a positive NAO is associated with strongly reduced blocking over Greenland and the northern North Atlantic, as in Shabbar *et al.* (2001); Croci-Maspoli *et al.* (2007) and Woollings *et al.* (2008). This is also consistent with the reduction in cyclonic wave-breaking seen in Figure 9. The increase in blocking to the south of the jet and over western Europe is also consistent with previous studies (Davini *et al.* 2013). Essentially the jet shifts southward due to blocking on its northern flank and northward due to blocking on its southern flank.

On the multidecadal timescale, blocking anomalies are less strongly related to the NAO, with only very weak anomalies at high and low latitudes. The implication is that the effect of blocking is largely to shift the jet stream whereas transient wave breaking can act both to shift or strengthen the jet depending on its position. Interestingly, there is an increase in blocking at the jet exit over the British Isles, despite the strengthening of the westerly winds there under the positive NAO. This may be a consequence of the storm track changes, since a strong storm track upstream is favourable for block maintenance (Shutts 1983; Nakamura and Wallace 1993). Häkkinen *et al.* (2011) demonstrated multidecadal variability of Atlantic-European blocking associated with Atlantic Ocean variability. Such basin-wide variations in blocking do not appear in the analysis presented here, suggesting that the NAO is not a good description of that variability.

7 Ocean-atmosphere interaction

The distinct physical characteristics of decadal NAO variability suggest an influence external to the atmosphere. The slowly varying ocean circulation is one potential forcing and a natural candidate is the Atlantic variability described by the Atlantic Multidecadal Oscillation (Knight *et al.* 2005; Sutton and Dong 2012). Correlations of the low-frequency NAO and jet indices with a smoothed AMO index are given in Figure 1. The AMO index was obtained from the NOAA ESRL website and was derived as in Enfield *et al.* (2001), including detrending and smoothing with a 121 month smoother. Annual means are plotted in Figure 1. The correlations show that the NAO is weakly anti-correlated with the AMO and that most of this correlation likely comes from the decadal variability in jet speed, which gives a slightly higher correlation of -0.48. Another potential candidate for ocean forcing of decadal NAO variability is the slow evolution of ocean temperatures in the tropical western Pacific (Kucharski *et al.* 2006; Manganello 2008).

These potential links are investigated in Figure 11 by correlating winter mean sea surface temperatures (SSTs) with the NAO at the two timescales. The SST data comes from the HadISST dataset (Rayner *et al.* 2003) and the correlation uses the ensemble mean NAO from the complete period of the 20CR data. Only gridpoints where the correlation is significant at the 95% level are shown. On the short timescale the SSTs show the familiar tripole pattern of anomalies which is largely a response to NAO variability. On the longer timescale the SSTs show a more global pattern, with significant values

outside of the Atlantic basin. The North Atlantic is generally cool, as expected from the negative correlation with the AMO, but the pattern is noisy and the large values elsewhere are hard to interpret and may not be physically related. The tropical western Pacific does, however, show a perturbed meridional SST gradient, as found by Kucharski *et al.* (2006).

Taking a similar approach in correlating the SSTs with the time series of jet speed from 20CR gives a clearer pattern on the multidecadal timescale, comprising a cold subpolar gyre in the North Atlantic and warm anomalies elsewhere, largely confined to the southern hemisphere (lower panels of Figure 11). Both of these features are reminiscent of AMO behaviour, and lends support to the potential role of Atlantic Ocean circulation in influencing NAO changes on the multidecadal timescale (e.g. Omrani *et al.* 2014; Peings and Magnusdottir 2014). There is also indication of a potential influence from the tropical Indian Ocean as suggested by Bader and Latif (2003).

In order to provide further evidence of an ocean influence, we examine the ocean-atmosphere coupling in more detail in the HiGEM simulation, where data availability and quality is not a limiting factor. Hodson and Sutton (2012) previously investigated the North Atlantic ocean-atmosphere coupling in this model with a focus on the shorter timescale. Figure 12 shows the correlation of winter (DJF) mean SST with the long timescale NAO for the model. This shows a distinct cold North Atlantic subpolar gyre; a pattern which is at least qualitatively similar to that related to jet speed in the observations (Figure 11). Figure 13a shows time series of the subpolar gyre temperature and the long

timescale NAO variability, which show reasonable covariability on the long timescale.

In order to determine an influence of the ocean on the atmosphere, Figure 13b presents the heat content budget for the subpolar gyre region. The sum of the heat contributions due to individual fluxes is shown by the dotted black line, and this agrees well with the total heat content in the thick black line, showing that all terms have been accounted for. The budget shows that it is the ocean heat flux convergence (in red) which is driving the heat content changes of the subpolar gyre, with the atmospheric fluxes (latent, sensible heating) acting to damp the changes in heat content. This agrees with the observational study of Gulev *et al.* (2013) which shows that ocean-atmosphere surface heat fluxes are driven by the ocean on multidecadal timescales.

The variations in ocean heat flux convergence into the subpolar gyre may be driven by a number of factors. Figure 13c demonstrates that these variations are closely related to variations in the Meridional Overturning Circulation at 45°N, suggesting that variations in meridional ocean transport are responsible. These variations in turn arise in response to west to east ocean pressure gradient across the Atlantic basin, which is dominated by ocean density variations on the deep western Atlantic boundary (Figure 13c: green line). The resulting picture is that the decadal variations in the SST of the subpolar gyre (see Figure 13a) are driven by changes in the MOC in the model, which are in turn driven by variations in the density within the deep western boundary current. Such density variations are ultimately generated at the ocean surface

in small regions of intense ocean cooling, such as the Labrador Sea (Marshall and Schott 1999). NAO variability is thought to be a significant factor in driving variations in ocean cooling in these regions (Eden and Willebrand 2001). Density anomalies generated by this process then slowly propagate southwards along the western Atlantic boundary.

In summary, subpolar SST anomalies in HiGEM arise due to changes in ocean heat convergence. These subpolar SST anomalies then in turn influence the atmosphere, likely by changing the meridional temperature gradient and hence the baroclinicity across the storm track. With a cold subpolar gyre the meridional gradient is strengthened which is expected to lead to stronger storm activity as seen in Figures 6 and 8. This in turn leads to increased acceleration of the westerly flow (Figure 8) and a stronger jet. Evidence of this mechanism of ocean influence on the atmosphere has been found in the natural variability of other models (Gastineau and Frankignoul 2012), in the context of model biases (Keeley *et al.* 2012) and in the response of models to climate change (Woollings *et al.* 2012).

8 Implications for long climate reconstructions

In this section we analyse the Küttel *et al.* (2010) reconstruction of Atlantic/European MSLP back to 1750 for evidence of contrasting NAO behaviour on the short and long timescales. The NAO in the reconstructed data is defined as the first EOF of winter mean MSLP over the region (0-40 °W, 20-70 °N), which roughly comprises the North Atlantic portion of the data domain. The surface

geostrophic zonal wind u_g is then derived from the MSLP and this is regressed onto the NAO at the two different timescales. The results are shown in the top panels of Figure 14. In contrast to the other datasets, the differences between the two regression patterns are small, with the NAO largely describing a jet shift on both timescales. To test whether this is due to the use of u_g rather than u_{850} , we apply the same procedure to derive u_g from 20CR and the results are shown in the middle panels of Figure 14. The results for 20CR resemble the difference in u_{850} found between a jet shift on short timescales and an increase in speed on longer timescales (Figure 3), suggesting that u_g from the reconstruction should be capable of capturing this behaviour.

The reconstruction method used in Küttel *et al.* (2010) is the multivariate principal component regression technique which relies on the assumption of stationarity. EOFs of both observed MSLP fields and pressure-sensitive proxy data (e.g., early measurements and documentary data such as ship log books) are combined with a multiple linear regression technique for the observational period to project local proxy information onto regional patterns. The linear relation is then assumed to be stationary over time and used to reconstruct MSLP fields further back in time. It is possible that the similarity of the NAO regressions on the two timescales is a consequence of this assumption of stationarity. To investigate this possibility we have performed a simple test on the 20CR data, treating it in an analogous way to the reconstruction method.

Firstly an EOF analysis of the MSLP is performed on the last 30 years of the 20CR data. Only the first four EOFs are retained, which explain 91% of

the variance of this sample. These EOFs are then used as a basis to truncate the full 136 year dataset: a multiple linear regression technique considering the anomaly maps for each year, the four leading EOFs and the associated principal component time series are used to derive MSLP fields by only including the projection on these four EOFs. The resulting pseudo-reconstructed MSLP fields are then regressed on the NAO at the two timescales and finally u_g is calculated from these. The results, shown in the lower panels of Figure 14, indicate some differences between the two timescales but these are substantially weaker than those in the original 20CR analysis (middle panels of Figure 14). This is particularly clear over Europe where the latitude of the wind anomalies is quite different in the middle panels but not in the lower panels. Retaining more than four EOFs (e.g. 10 EOFs which comprises 99% of the variance) does not significantly alter these findings (not shown). While this procedure is analogous but not identical to the technique of Küttel *et al.* (2010), it does suggest that the lack of a distinct multidecadal NAO signature in the reconstruction could be at least partly due to the assumption of stationarity in the method. This effect could be compounded by the relatively low density of proxy data in the jet stream region over the ocean considered in the reconstruction of Küttel *et al.* (2010), and non-climatic noise intrinsic to proxy data.

9 Conclusions

This study shows that the multidecadal variability of the NAO represents very different variations in atmospheric circulation from the interannual-decadal

variability. The faster variability is dominated by meridional shifts of the jet stream and associated storm track, while the slower variability is dominated by changes in the speed of the jet and the strength of the storm track.

Variations on both timescales are supported by forcing from the transient eddies, but the nature of this forcing is different. The interannual-decadal variations are associated with shifts of the transient eddy forcing and with the occurrence of blocking weather patterns. Other work suggests this variability represents variations in the occurrence of different synoptic circulation regimes such as preferred jet positions (W14 and references therein). In contrast, the multidecadal variability is associated with changes in strength of the eddy forcing and with in-phase changes in the occurrence of transient Rossby wave breaking on both sides of the jet.

The patterns of influence of the NAO on regional temperatures, wind speeds and precipitation are different on the two timescales, and this has clear implications for the interpretation of proxy or reconstructed records of past atmospheric variability in this region. The variations on multidecadal timescales may not be well represented by the canonical NAO pattern, especially since the shorter timescale variability dominates the variance of the NAO index. A potential example of this has been given, by analysing a long MSLP reconstruction. In contrast to the other datasets, this does not exhibit a difference in NAO character on short and long timescales, and it is suggested that the stationarity assumption commonly used in reconstruction methods is at least

partly responsible for this. These findings have implications for the interpretation of climate reconstructions and long climate records.

These results also provide strong evidence for the presence of some forcing on the decadal NAO from more slowly varying components of the climate system than the atmosphere. Some evidence of links to Atlantic Ocean variability were revealed, although other factors may also contribute. This evidence is particularly clear in the HiGEM GCM, where variations in the Atlantic Meridional Overturning Circulation lead to significant SST anomalies in the subpolar gyre region which are then damped by the heat fluxes to the atmosphere.

For the emerging discipline of decadal prediction these results are an encouraging sign of potential predictability of the winter NAO on multidecadal timescales. Furthermore, the multidecadal component of NAO variability has a clear and distinct influence on surface temperatures and precipitation, especially in Europe, so that decadal forecasts of this variability could be of practical use. However, the contrasting behaviour on interannual and decadal timescales suggests that the potential sources of skill may be different for decadal forecasts than for seasonal forecasts.

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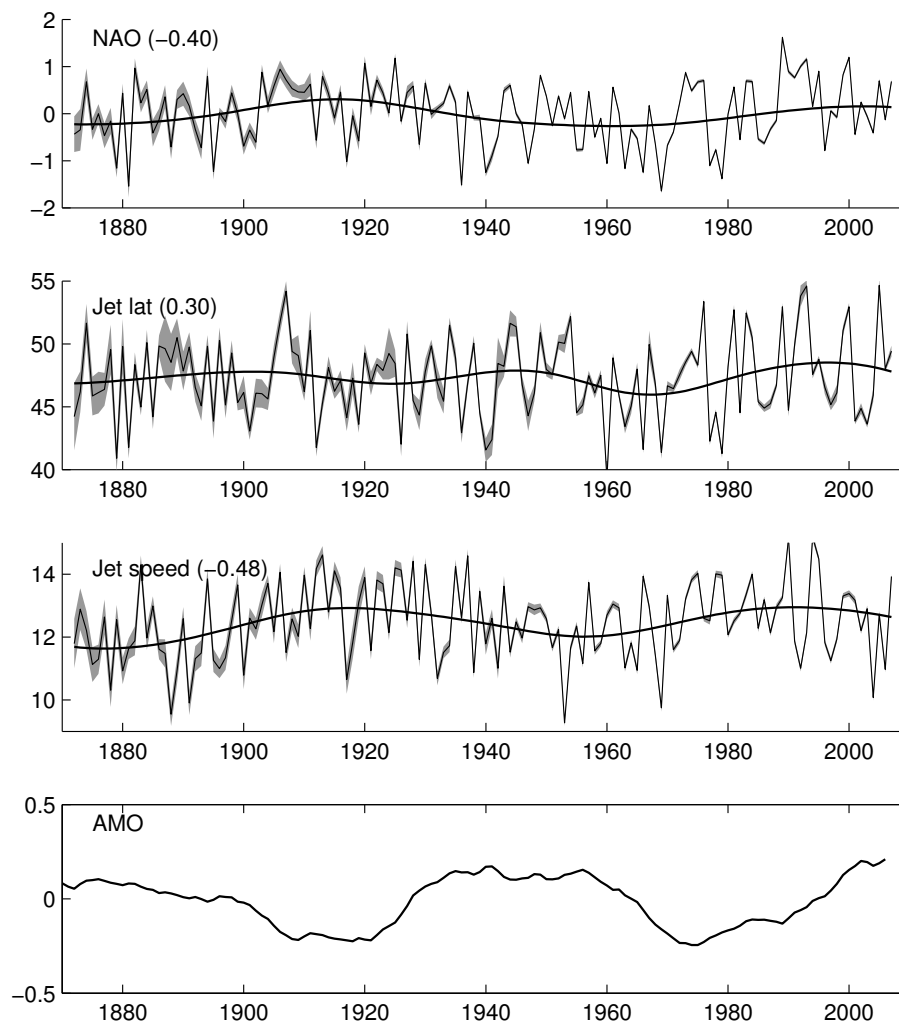


Fig. 1 Ensemble mean indices of the winter mean NAO, jet latitude and jet speed from 20CR, with the multidecadal (>30 year) component also shown. The shading indicates the ± 2 standard deviation range across the ensemble. The AMO is shown in the bottom panel, and in each other panel the correlation of the respective low-frequency timescale with this is given.

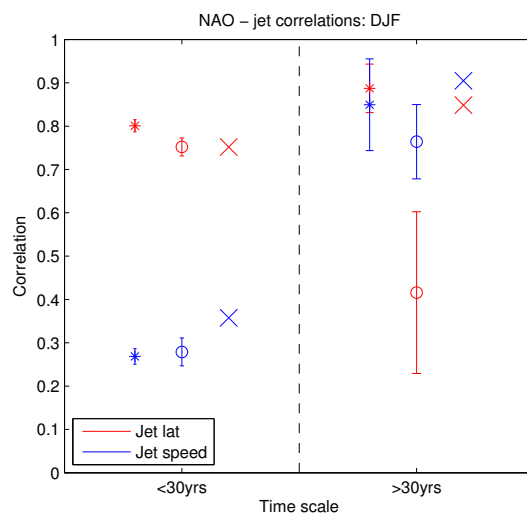


Fig. 2 Correlations of jet indices with the NAO on short and long timescales, showing NCEP-NCAR results as crosses and 20CR as circles, with errorbars giving the ± 2 standard deviation range across the ensemble. Asterisks give 20CR results for the NCEP-NCAR period.

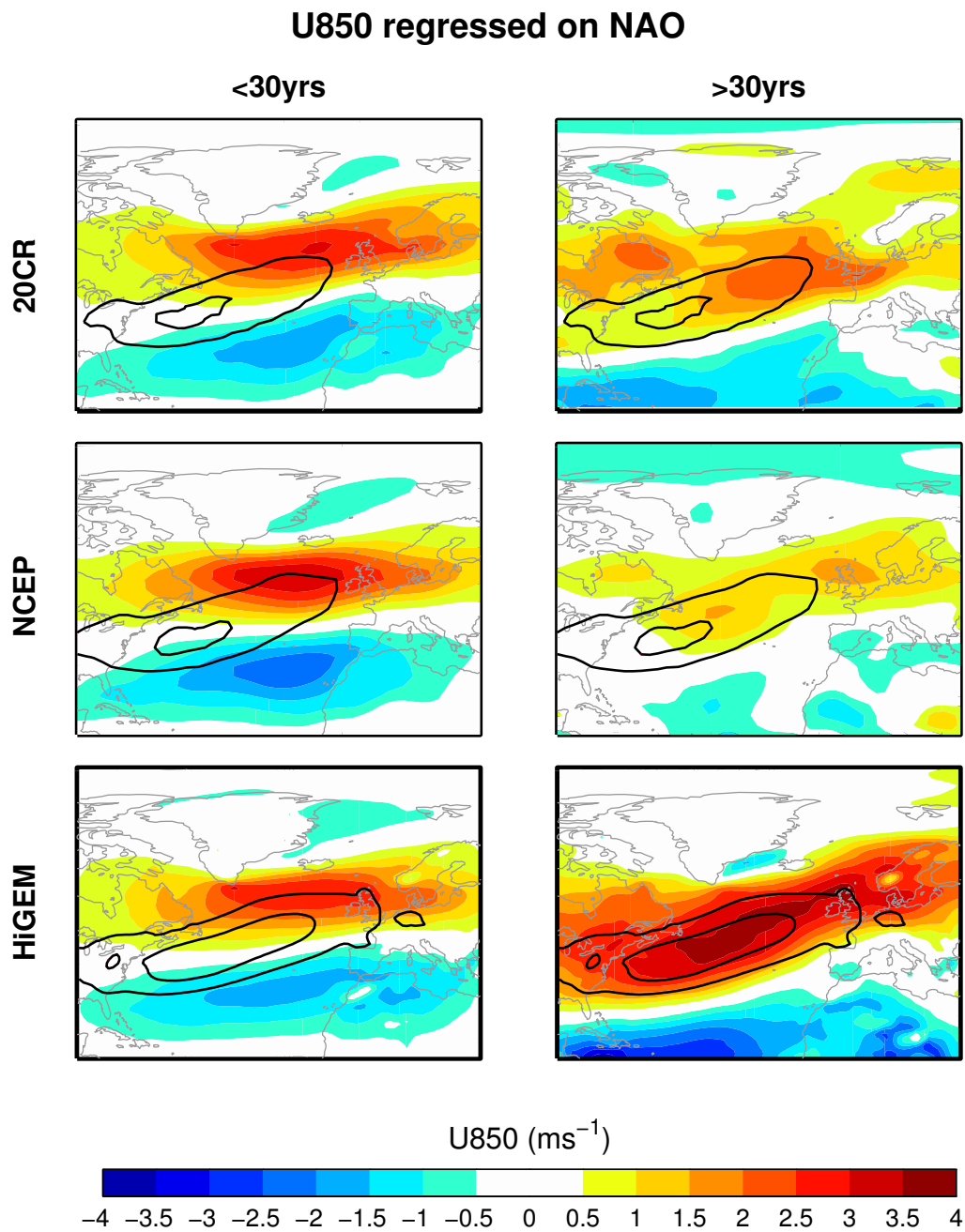


Fig. 3 Regression patterns of anomalies in 850hPa zonal wind on the NAO at the two timescales, using 20CR, NCEP-NCAR and HiGEM. The wind climatology is shown in black contours at 7.5 and 10.5 ms^{-1} .

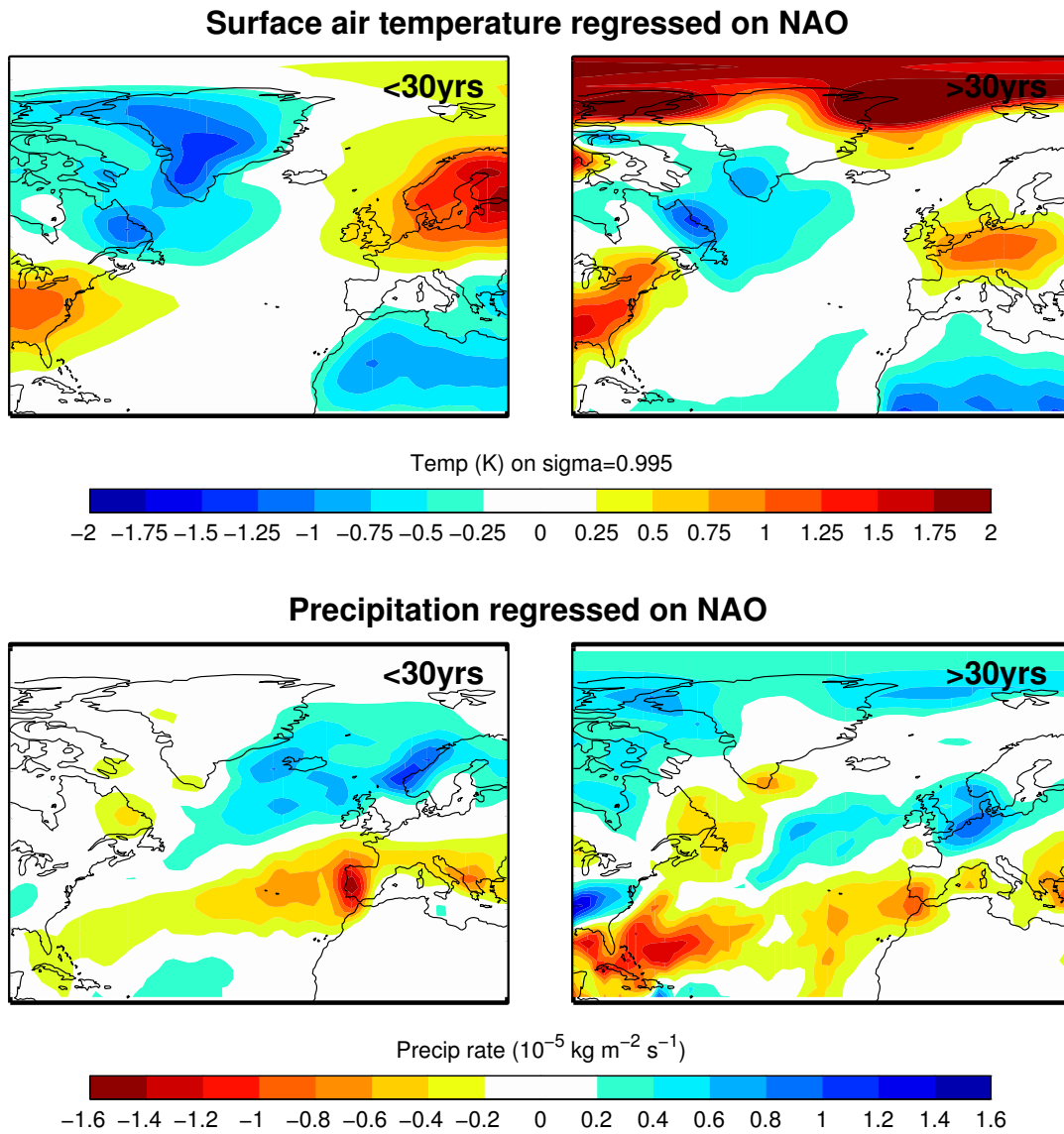


Fig. 4 Regression patterns of near surface air temperature (on the $\sigma = 0.995$ level) and precipitation on the NAO at the two timescales. This analysis was performed on the ensemble mean fields from 20CR.

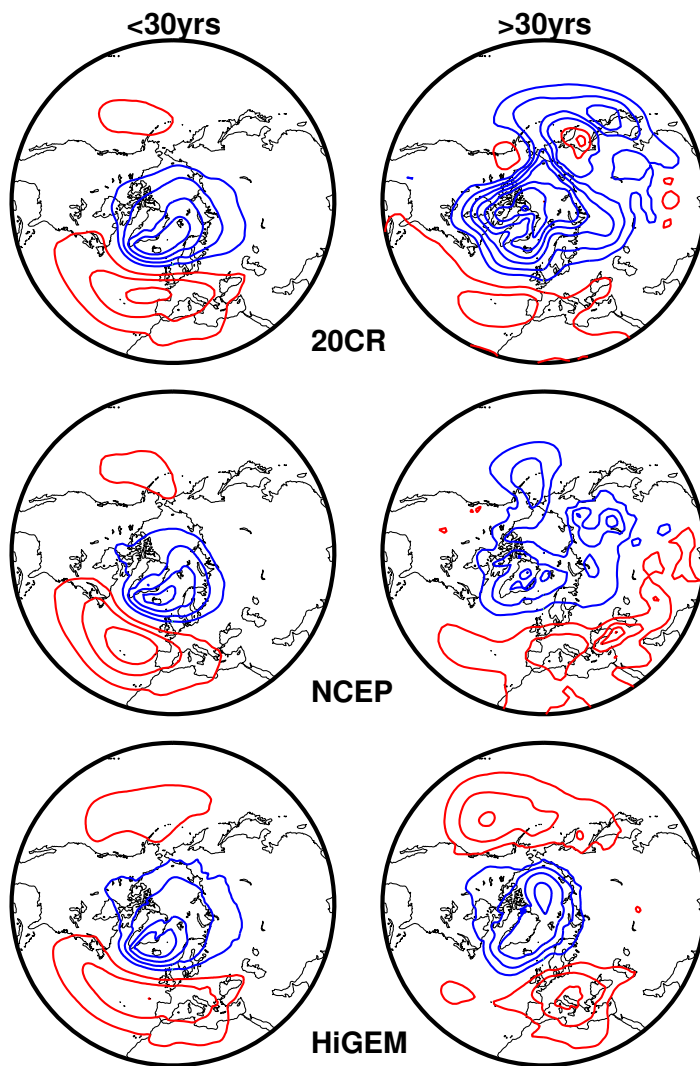


Fig. 5 MSLP from both reanalyses and from the HiGEM model regressed onto the NAO on both timescales. Contours are drawn every 1 hPa with negative contours in blue. All ensemble members are used for 20CR and the results averaged.

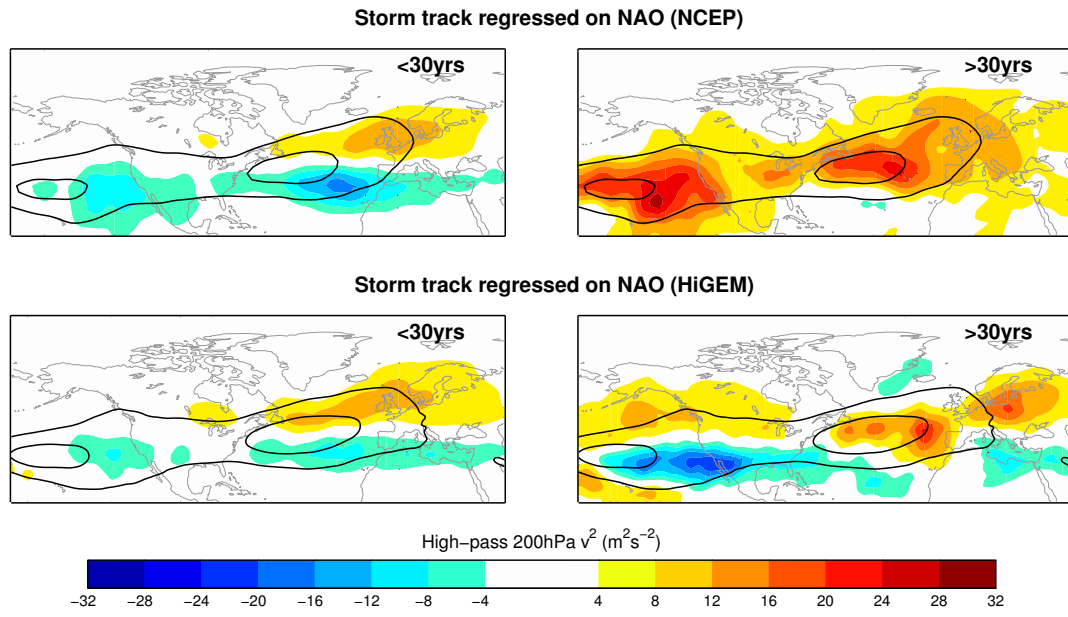


Fig. 6 Regressions of the storm track activity on the NAO in the NCEP-NCAR reanalysis and HiGEM, using the square of the meridional wind anomalies after applying a 2-6 day filter. The climatology is contoured at 60 and 100 m^2s^{-2} .

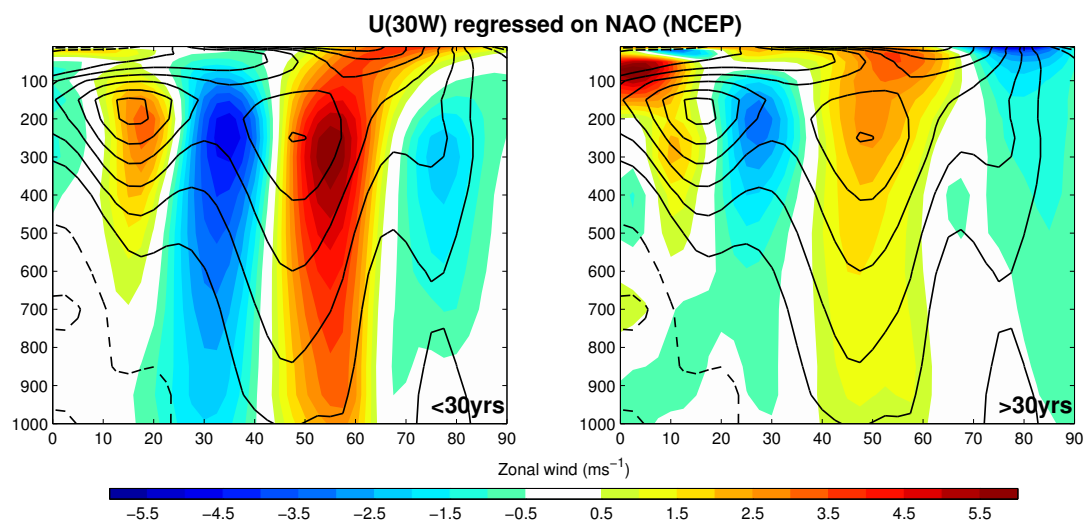


Fig. 7 Regression patterns of anomalies in zonal wind at 30°W on the NAO at the interannual and decadal timescales, using the NCEP reanalysis. The wind climatology is contoured every 5 ms^{-1} .

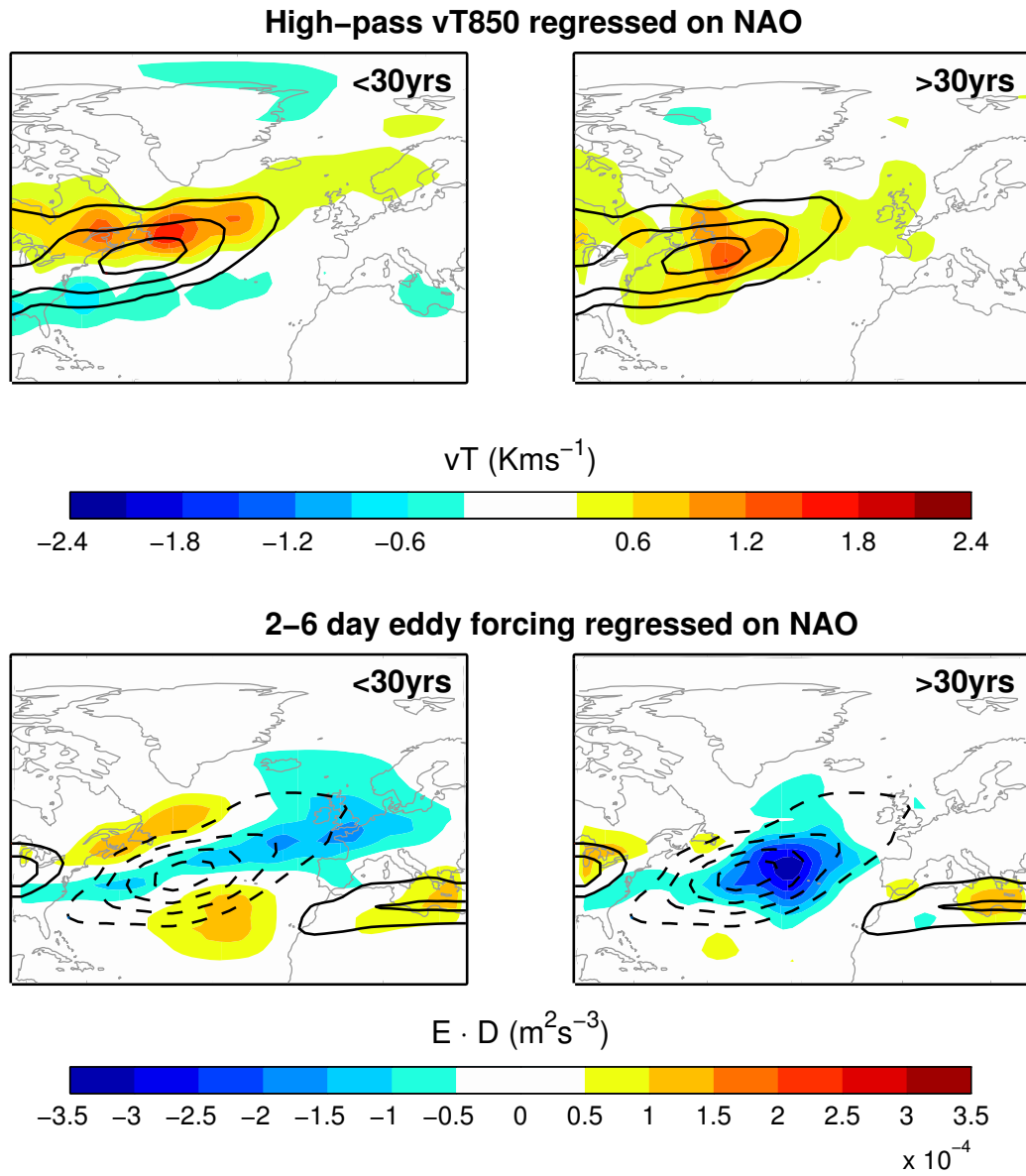


Fig. 8 Top: Regressions on the NAO of the 2–6 day $v'T'$ at 850 hPa. The climatology is contoured at 4, 7 and 10 Kms^{-1} . Bottom: Regressions on the NAO of the eddy forcing diagnostic $E \cdot D$ at 250 hPa. The climatology is contoured every $5 \text{ m}^2\text{s}^{-3}$, with negative contours dashed and the zero contour omitted. In all cases the data is from the NCEP–NCAR reanalysis.

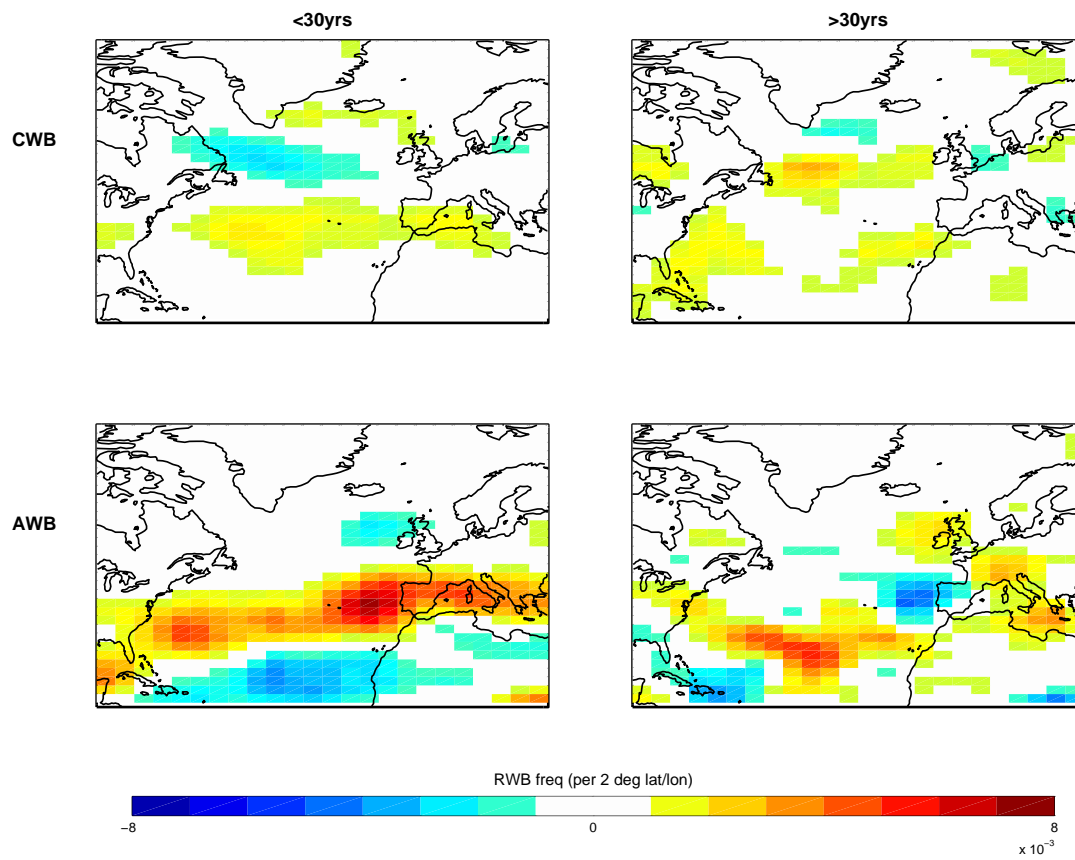


Fig. 9 Regressions on the NAO of the transient Rossby wave breaking occurrence, split into cyclonic (CWB) and anticyclonic (AWB). The NCEP-NCAR reanalysis is used.

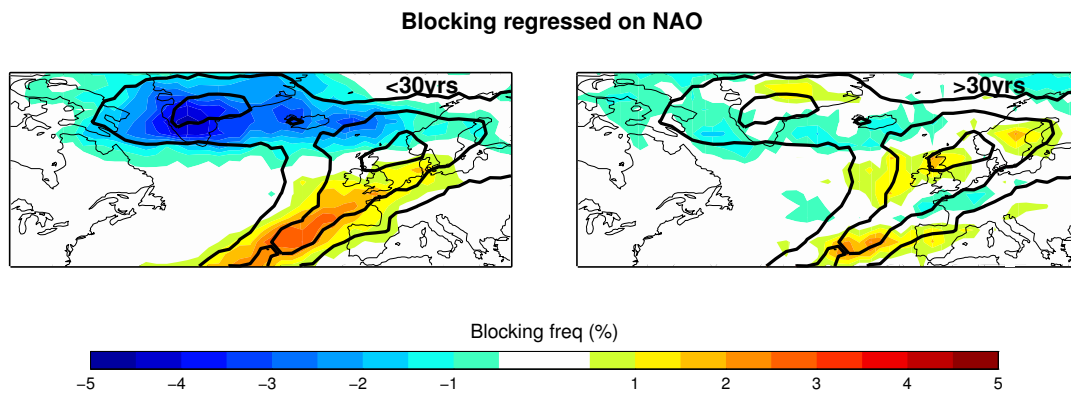


Fig. 10 Regressions of blocking occurrence on the NAO, using the NCEP-NCAR reanalysis.

The climatology is shown in black contours every 2%.

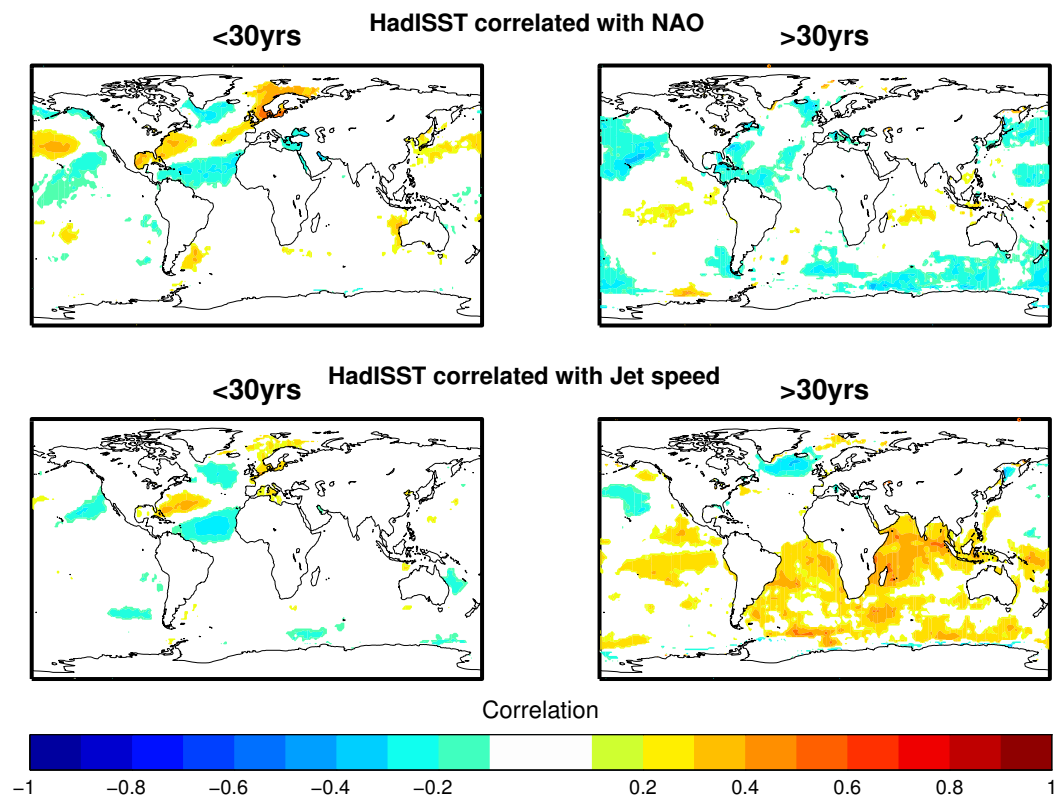


Fig. 11 Correlation value r of winter mean SSTs on the NAO (top) and jet speed (bottom) at the two timescales. Only values which are significant at the 95% level have been plotted.

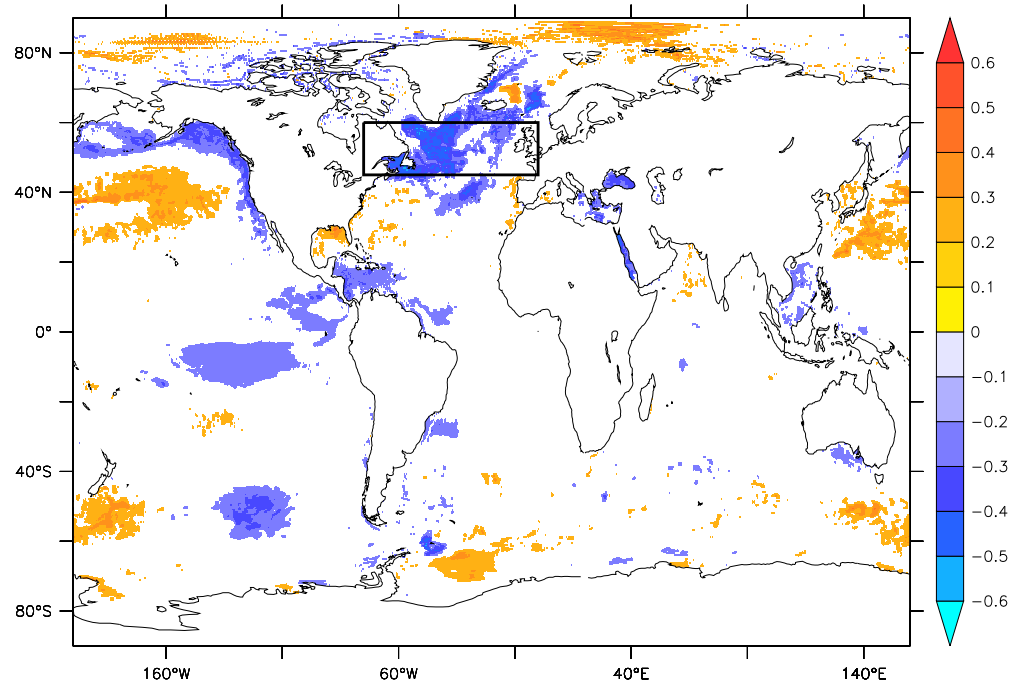


Fig. 12 Correlation value r of Sea Surface Temperature (SST) correlated with multidecadal NAO index in HiGEM. Shaded areas are significant at the 95% level ($p < 0.05$). Both SST and NAO were detrended before correlation.

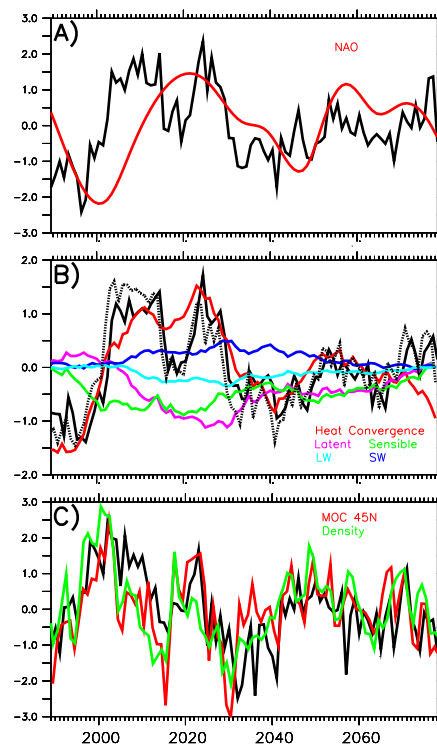


Fig. 13 Decadal NAO variability and Sub Polar Gyre heating in HiGEM Control simulation. A) Black: Mean Atlantic Sub Polar Gyre (SPG) Sea Surface Temperature (SST) (75:0°W, 45:60°N- box in Figure 12). Red: Detrended decadal component of the NAO in HiGEM multiplied by -1, extracted using EMD as for observations. Both indices have been standardized to have unit variance. B) Heat budget for the SPG region. Black solid: upper ocean heat content within the SPG region (0:500 m depth). Other lines - Heat content in the SPG due to: Ocean Heat convergence (Red), Surface Latent (Purple) and Sensible (Green) Heat fluxes and Longwave (Light Blue) and Shortwave (Dark Blue) surface radiation fluxes. All surface fluxes are defined positive *into* the ocean. Black dotted line: the sum of all contributions to the heat content. All indices have been detrended. Units are 10^7 PJ. Black (solid and dotted) lines have been multiplied by 2 to aid comparison with SPG SST in panel A. C) Black: Ocean Heat Convergence Flux into the SPG region (45:60°N). Red: Atlantic Meridional Overturning Circulation (AMOC) at 45°N (AMOC is the integral of southward meridional ocean velocity between 1000:7000 m across the Atlantic Basin). Green: Mean Ocean Density on the Deep western Atlantic Boundary (1500:3000 m 59:58°W 44:45°N). All indices have been detrended and standardized to have unit variance.

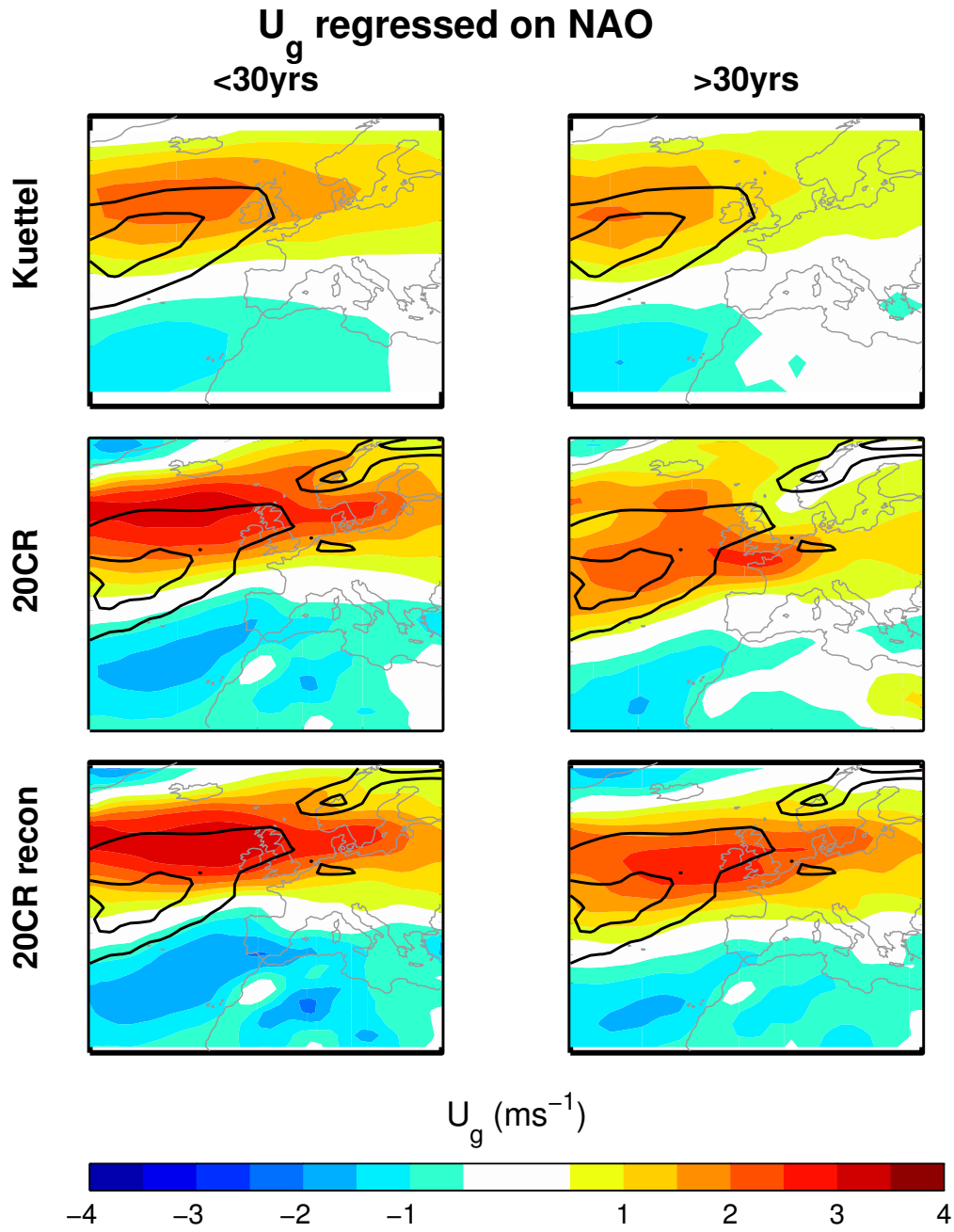


Fig. 14 As Figure 3 but showing the surface geostrophic zonal wind using the Küttel reconstruction and the 20CR data. The wind climatology is contoured in black at $\pm 5, 7 \text{ ms}^{-1}$. The reconstruction only covers the region shown. See text for further details.